

# Radial Conformal Lattice Field Theory: Spherical Manifolds

$$\mathbb{R}^d \rightarrow \mathbb{R} \times \mathbb{S}^{d-1}$$

*Richard Brower Boston University*

*Lattice 2014 June 25*

\*RCB, G. Fleming, H. Neuberger, M. Cheng

# Radial Quantization: *Early History*

- ▣ S. Fubini, A. Hanson and R. Jackiw PRD 7, 1732 (1972)

**Abstract:** A field theory is quantized covariantly on Lorentz-invariant surfaces. Dilatations replace time translations as dynamical equations of motion. .... The Virasoro algebra of the dual resonance model is derived in a wide class of 2-dimensional Euclidean field theories.

- ▣ J. Cardy J. Math. Gen 18 757 (1985).

**Abstract:** The relationship between the correlation length and critical exponents in finite width strips in two dimensions is generalised to cylindrical geometries of arbitrary dimensionality  $d$ . For  $d > 2$  these correspond however, to curved spaces. The result is verified for the spherical model

# Radial Quantization

Evolution:  $H = P_0$  in  $t \implies D$  in  $\tau = \log(r)$

$$ds^2 = dx^\mu dx_\mu = e^{2\tau} [d\tau^2 + d\Omega^2]$$

Can drop  
Weyl factor!

$$\mathbb{R}^d \rightarrow \mathbb{R} \times \mathbb{S}^{d-1}$$

"time"  $\tau = \log(r)$ , "mass"  $\Delta = d/2 - 1 + \eta$

$$D \rightarrow x_\mu \partial_\mu = r \partial_r = \frac{\partial}{\partial \tau}$$

# Motivation

- (near) Conformal Field Theories, interesting for
  - BSM composite Higgs
  - AdS/CFT weak-strong duality
  - Model building & Critical Phenomena in general
- Lattices on  $\mathbb{R}^d$  are problematic:
  - scales are exponentially divergent.
- Linear Hypercubic vs Exponential Radial Lattice

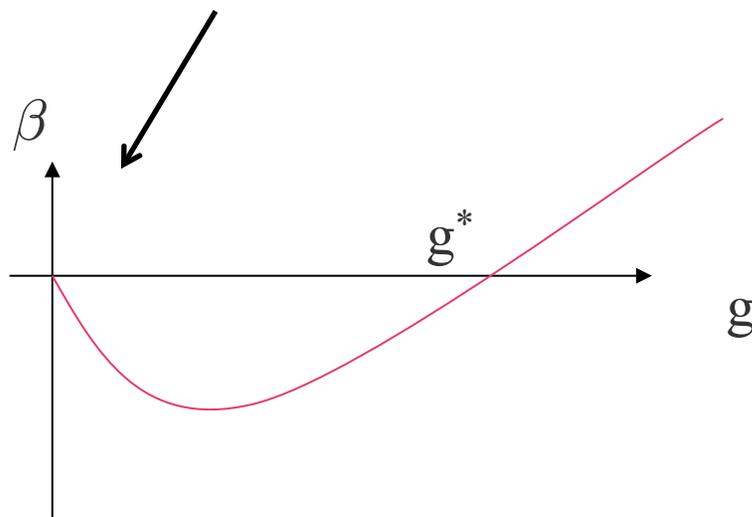
$$a < \Delta r < L \quad \text{vs} \quad a < \Delta \log(r) < L$$

An IR fixed point can emerge already in the two-loop  $\beta$ -function as you increase the number  $N_f$  of fermions.

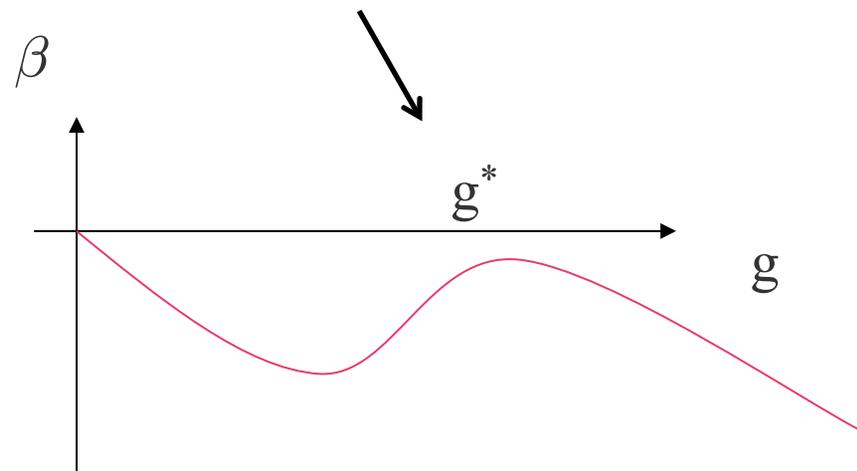
$$\beta(g) = b_0 g^3 + b_1 g^5 + \dots$$

$b_0 < 0$  for  $N_f < 11N_c/2 = 16.5$

$b_1 > 0$  for  $N_f > 153/14$



Conformal Window:  
 $16.5 > N_f > N_f^*$



Near conformal  
 $N_f < N_f^*$  , close to  $N_f^*$

# CFT are highly constrained

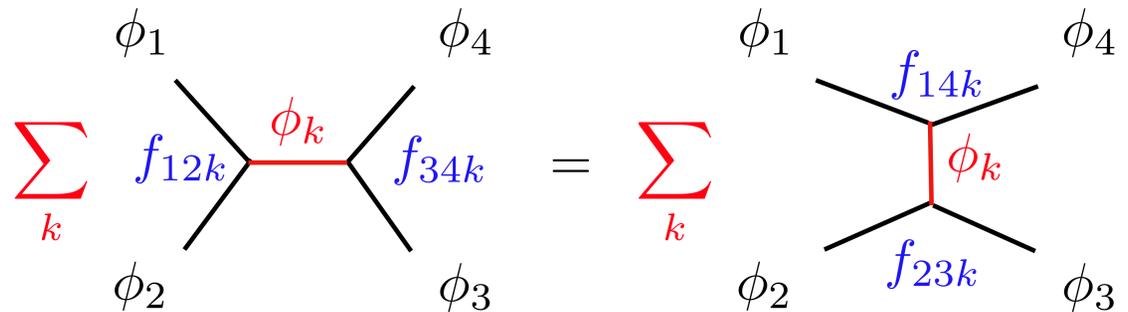
(i.e. Data: spectra + couplings to conformal blocks)

Exact 2 and 3  
correlators

$$\langle \phi(x_1)\phi(x_2) \rangle = C \frac{1}{|x_1 - x_2|^{2\Delta}}$$

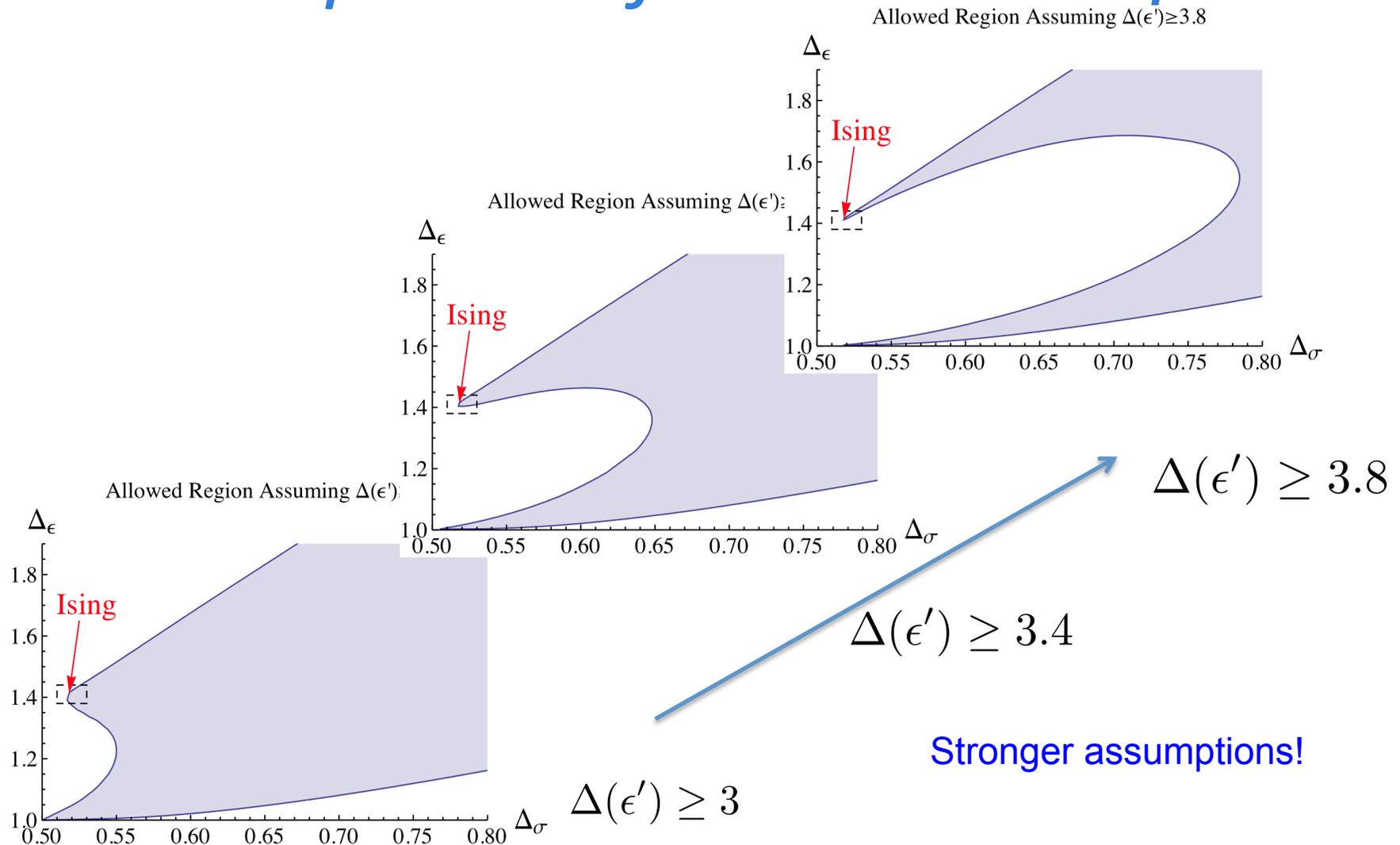
$$\mathcal{O}_i(x_1)\mathcal{O}_j(x_2) = \sum_k \frac{f_{ijk}}{|x_1 - x_2|^{\Delta_i + \Delta_j - \Delta_k}} \mathcal{O}_k(0)$$

Only "tree" diagrams!  
"partial waves" exp: sum  
over conformal blocks



CFT Bootstrap: OPE & factorization completely fixed the theory

# Inequalities from Bootstrap\*



•“Solving the 3D Ising Model with the Conformal Bootstrap” (El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin and Vichi) arXiv:1203.6064v1v [hep-th] (2012)

# *Narrative*

## I. First attempt

-- *Lattice Radial Quantization: 3D Ising*

R.C.B., G.T. Fleming and H. Neuberger, Phys. Lett. B 721 (2013)

## II. What worked and what failed.

## III. Finite Elements Methods (FEM) to the rescue ?

## IV. Future hopes and dreams

# Exact CFT: Power Law Correlator

Conformal correlator:  $\langle \phi(x_1)\phi(x_2) \rangle = C \frac{1}{|x_1 - x_2|^{2\Delta}}$

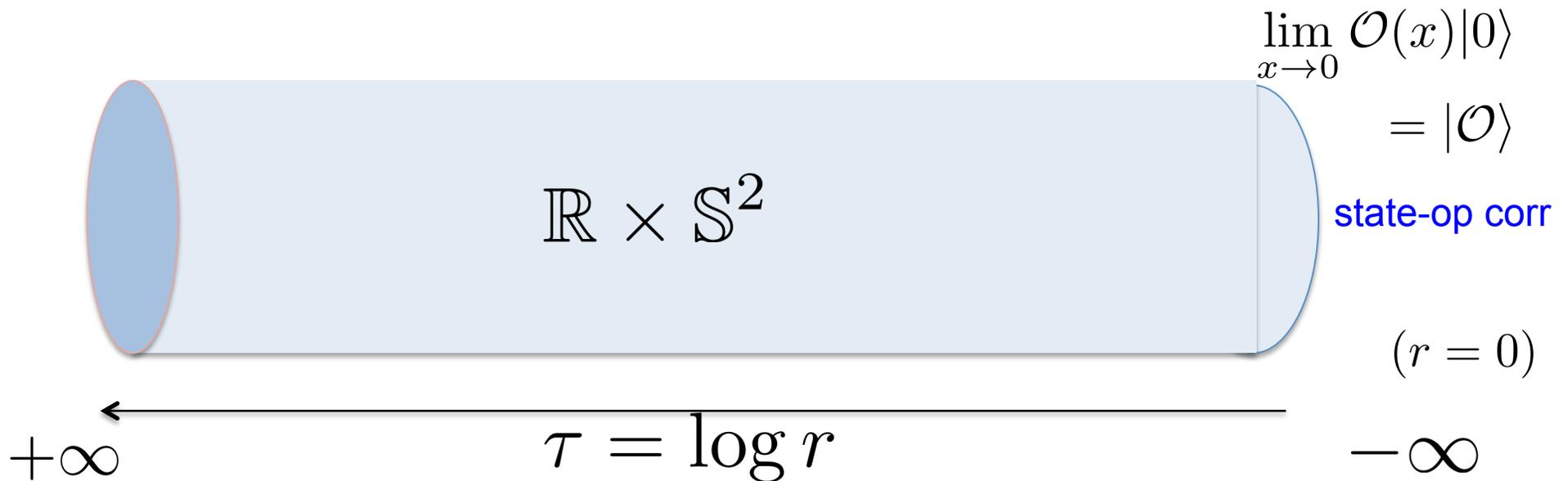
$$\begin{aligned} r_1^\Delta r_2^\Delta \langle \phi(\tau_1, \Omega_1)\phi(\tau_2, \Omega_2) \rangle &= C \frac{1}{[r_2/r_1 + r_1/r_2 - 2 \cos(\theta_{12})]^\Delta} \\ &\simeq C e^{-(\log(r_2) - \log(r_1))\Delta} \\ &= C e^{-\tau\Delta} \end{aligned}$$

With  $|x_1 - x_2|^2 = r_1 r_2 [r_2/r_1 + r_1/r_2 - 2 \cos(\theta_{12})]$

as  $\tau = \log(r_2) - \log(r_1) \rightarrow \infty$

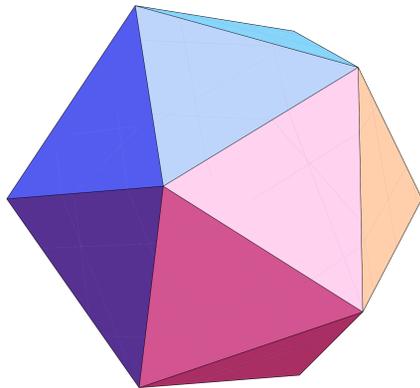
# 3-d Ising at Wilson-Fisher FP

$$Z_{Ising} = \sum_{\sigma(x,t)=\pm 1} e^{\beta \sum_{t, \langle x,y \rangle} \sigma(t,x)\sigma(t,y) + \beta \sum_{t,x} \sigma(t+1,x)\sigma(t,x)}$$

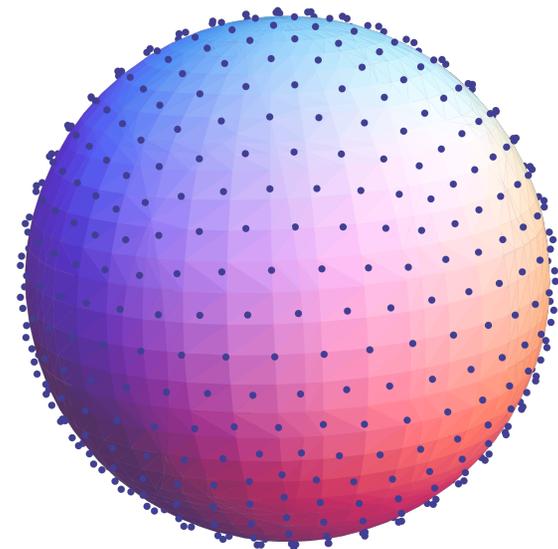
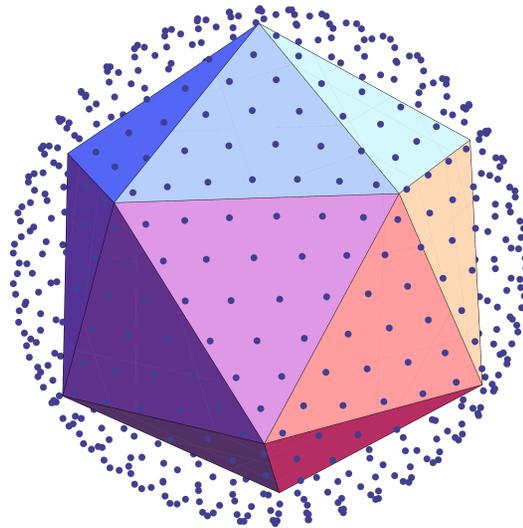


# *Order $s$ Refined Triangulated Icosahedron*

$s = 1$



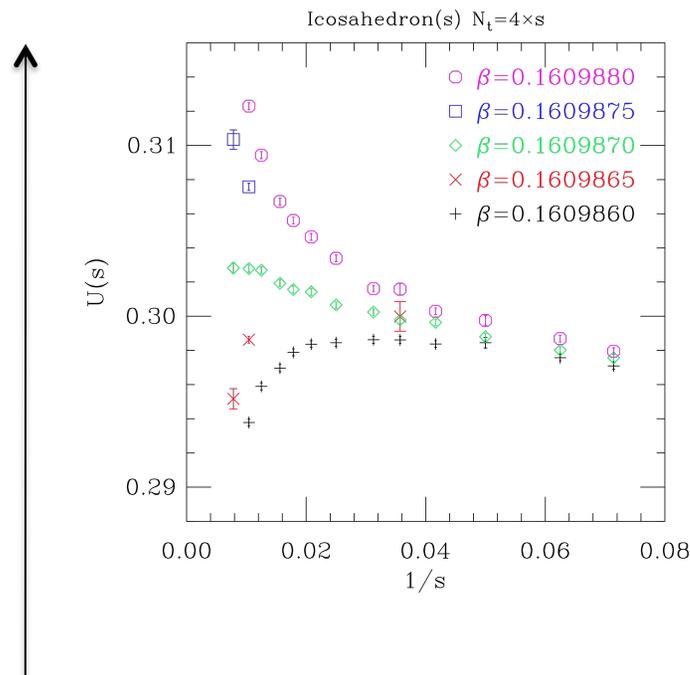
$s = 8$



$l = 0$  (A),  $1$  (T1),  $2$  (H) are irreducible 120  
Icosahedral subgroup of  $O(3)$

# Fitting to Finite scaling

$$U[(\beta - \beta_{cr})L^{1/\nu}, (\lambda - \lambda_{cr})L^{-\omega}, \dots] \simeq U^*(x) + O(L^{-\omega}) \simeq U^*(0) + a_1(\beta - \beta_{cr})L^{1/\nu} + c(\lambda)L^{-\omega} + \dots$$



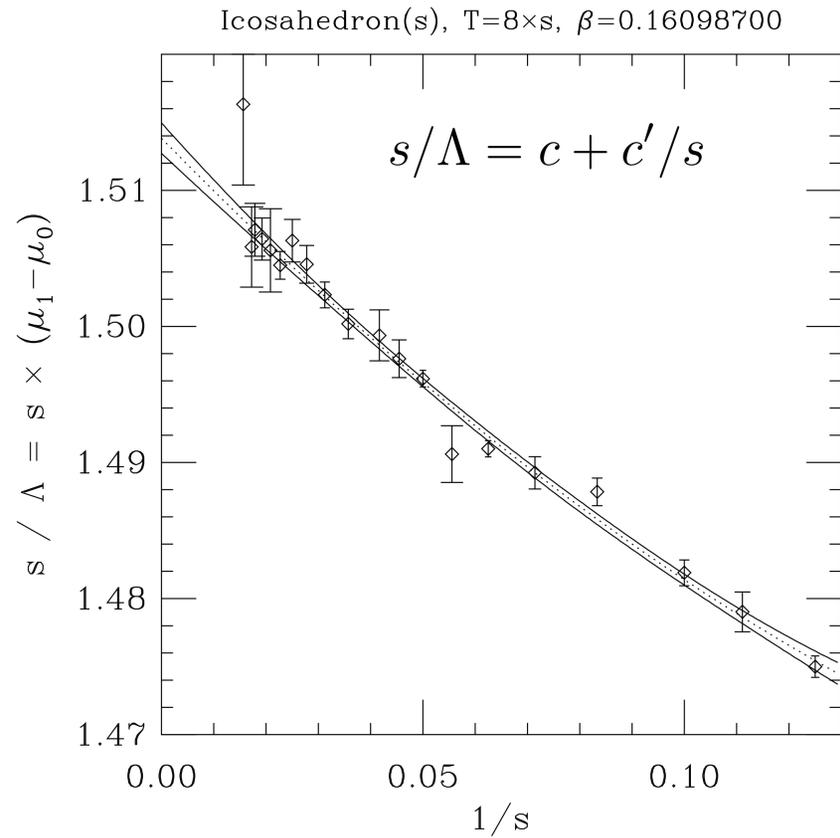
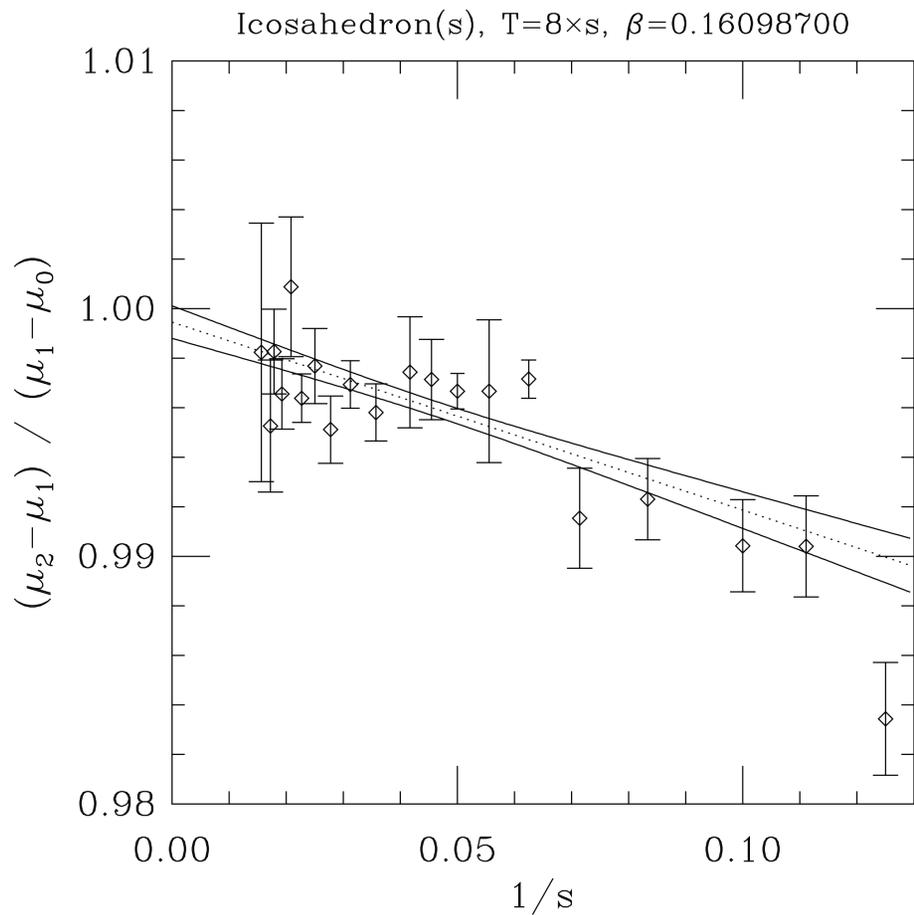
$$U(\beta, s = L) = 1 - \frac{\langle M^4 \rangle}{3\langle M^2 \rangle^2}$$

$$\beta_{cross} \simeq \beta_{cr} + c_1 L^{-1/\nu - \omega}$$

$$\beta_{crit} = 0.16098703(3)$$

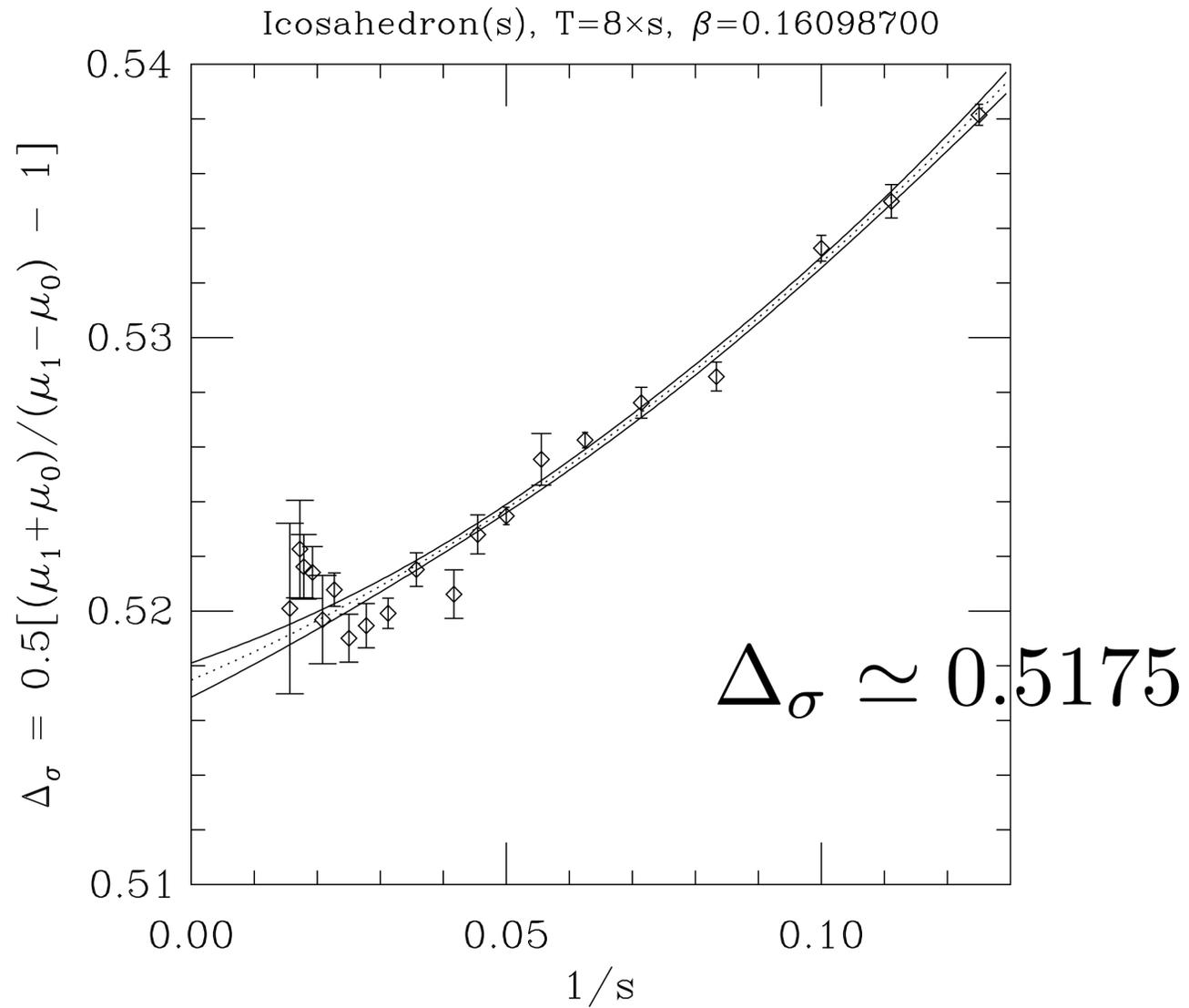
Double Scaling:  $x = (\beta - \beta_{cr})L^{1/\nu}$

*Determine “speed of light” via  
Descendant Relation & rescale “log(r)”*



$$c = 1.5105(7)$$

# Current Fit:



# Improved cluster Estimator

Swendsen-Wang: Real space

$$g(x - y) = \langle s_x s_y \rangle \simeq \frac{1}{N_{config}} \sum_{i=1}^{N_{config}} \sum_{C_i} \Delta_{C_i}(x) \Delta_{C_i}(y)$$
$$\Delta_C(x) = 1 \text{ if } x \in C \text{ else } 0$$

Wolff single cluster

$$\tilde{g}_{lm}(k) \simeq \frac{1}{N_{config}} \sum_{i=1}^{N_{config}} \frac{1}{|C|} \left| \sum_{t,x \in C} e^{i2\pi kt/L_t} Y_{lm}(\Omega_x) \right|^2$$

Note: All to All O(V) improved estimator in Momentum space \*

\*C. Ruge, P. Zhu and F. Wagner Physica A (1994) 431:

# Numerical Test

- Equal spacing test of descendants:

$$\frac{\mu_2 - \mu_1}{\mu_1 - \mu_0} = 0.999(1)$$

- “Speed of light”  $c = 1.5105(7)$

- But critical point  $\beta_{crit} = 0.16098703(3)$

- Current anomalous dimensions (more soon)

- from Binder:  $\omega + 1/\nu = 2.51(11)$
- from corr:  $\Delta_\sigma = 1/2 + \eta/2 = 0.5175(6)$
- Simulation are on going to reduce errors

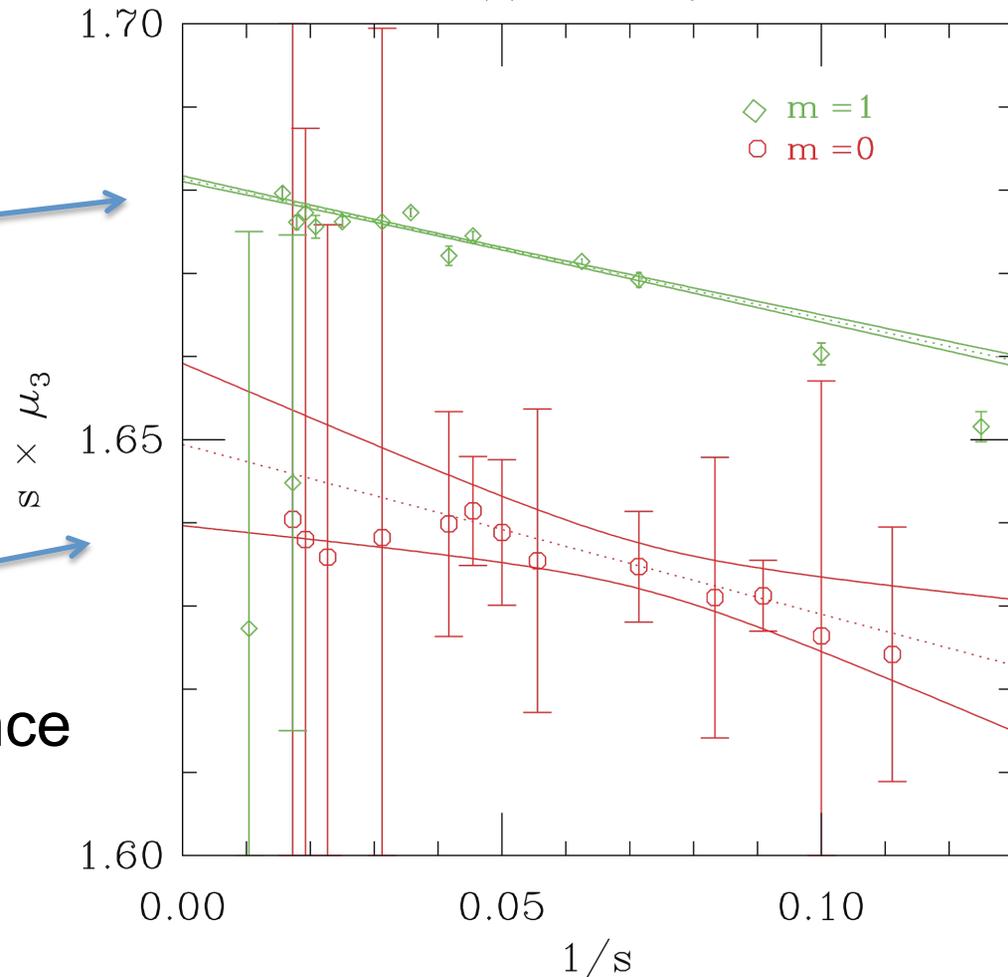
# Wrong Theory?

## Failure to recover $O(4,1)$ of $l = 3$ ?

Icosahedron(s),  $T=8 \times s$ ,  $\beta=0.16098700$

G rep

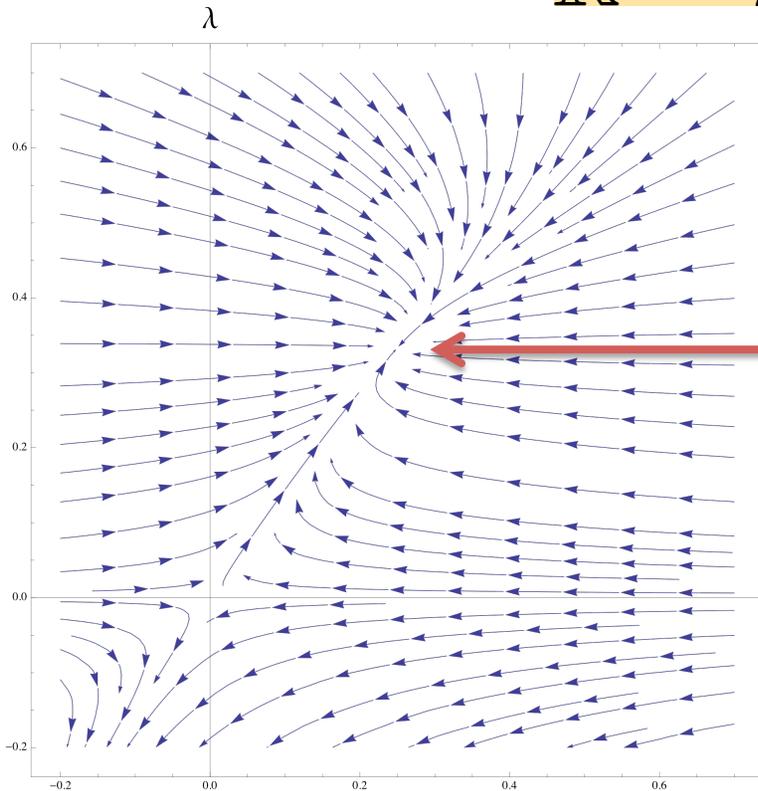
T2 rep



Apparent lack of convergence  
to a single  $O(3)$  irreducible  
representation for  $l = 3$

# Replace Ising Model by phi 4<sup>th</sup>

$$\mathbb{R}^d \rightarrow \mathbb{R} \times \mathbb{S}^{d-1}$$



$$\mathcal{L}(x) = -\frac{1}{2}(\nabla\phi)^2 + \lambda(\phi^2 - \mu/2\lambda)^2$$

Wilson-Fisher FP

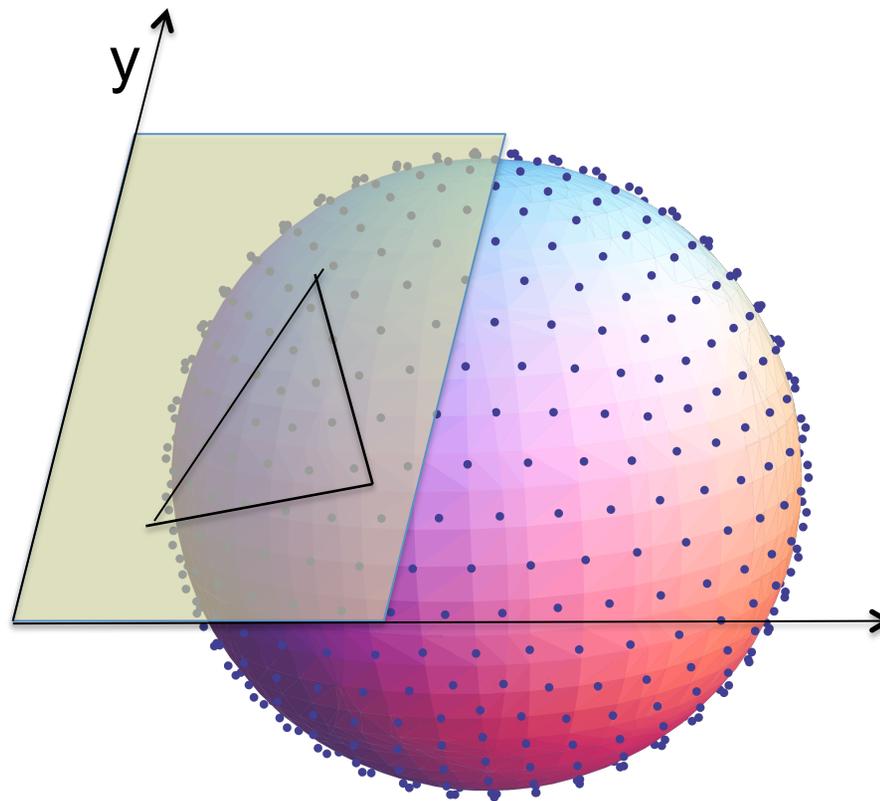
Gaussian FP

$$\beta_g = \epsilon g - \frac{3}{16\pi^2}g^2 + O(g^3, \epsilon g^2, \mu^4, \mu^2 g; )$$

$$\beta_{\mu^2} = 2\mu^2 + ag + \frac{9}{16\pi^2}g\mu^2 + O(\mu^4) \quad \lambda = 4g/4!$$

# *Finite element Discrete Lagrangian on triangulated sphere.*

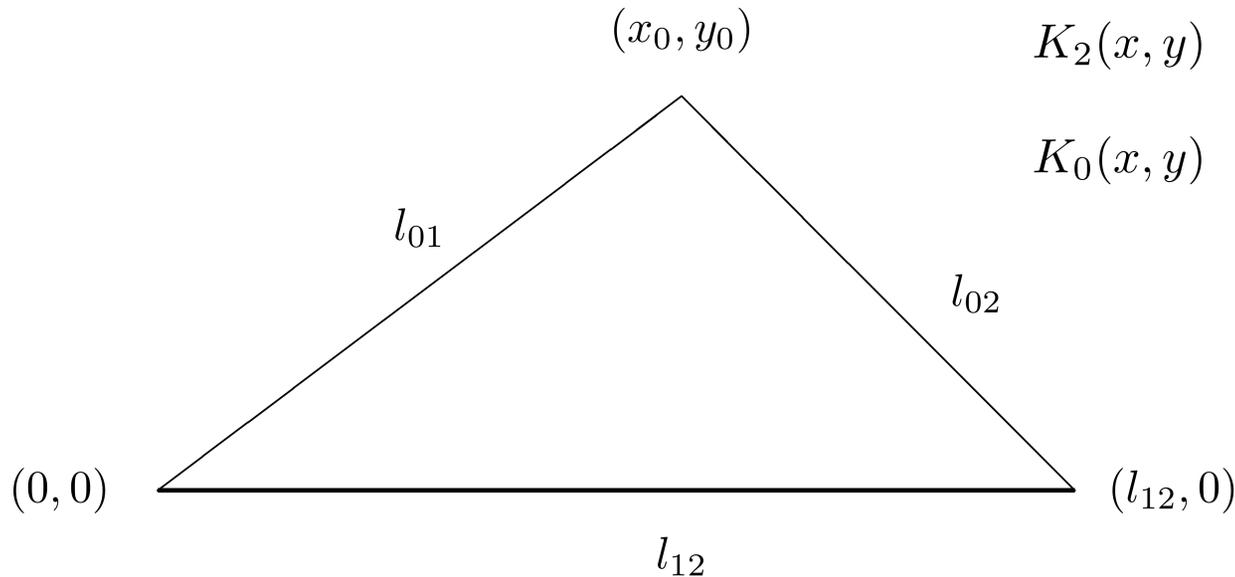
$$Z = \int \mathcal{D}\phi e^{-\frac{1}{2} K_{ij} (\phi_i - \phi_j)^2 - \lambda \omega_i (\phi^2 - \mu/2\lambda)^2}$$



project spherical  
triangle onto  
local tangent plane

x

# Kinetic term for Linear Element



$$K_1(x, y) = \left[ l_{12} - x - \frac{(l_{12} - x_0)y}{y_0} \right] / l_{12}$$

$$K_2(x, y) = \left[ x - \frac{x_0 y}{y_0} \right] / l_{12}$$

$$K_0(x, y) = \frac{y}{y_0}$$

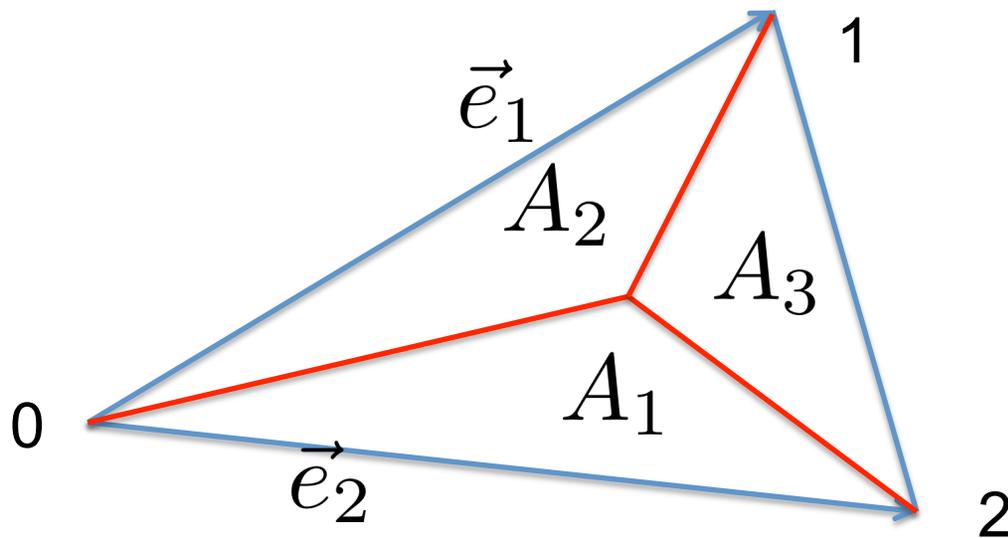
On each triangle expand:  $\phi(x, y) = \sum_i K_i(x, y) \phi_i$  and integrate

$$\int_{A_{012}} dx dy \partial_\mu \phi(x, y) \partial_\mu \phi(x, y) = \frac{1}{2A_{012}} [(l_{01}^2 + l_{20}^2 - l_{12}^2)(\phi_1 - \phi_2)^2 + \text{cyclic}]$$

See also: Christ, Friedberg, Lee on "Random Lattice" NP (1982)

## Linear Finite Element Method for triangulate Manifold

$$g_{ij}(0) = \vec{e}_i \cdot \vec{e}_j$$



$$K_i^\Delta(r) = \frac{A_i}{A_{123}}$$

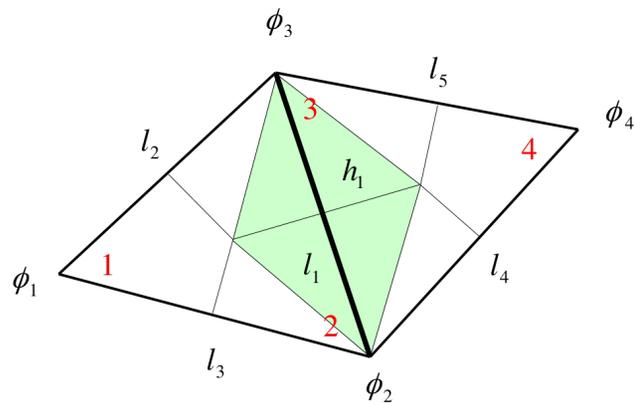
$$W_i(r) = \sum_{\Delta} K_i^\Delta(r)$$

$$\phi(r) = \sum_i W_i(r) \phi_i$$

Piecewise linear  
subspace of Hilbert space

$$d\vec{x} = \vec{e}_i dx^i \quad ds^2 = d\vec{x} \cdot d\vec{x} = g_{ij} dx^i dx^j$$

Regge Calculus formulation for smooth manifold.



$$FEM : A_d \frac{(\phi_2 - \phi_3)^2}{l_1^2}$$

Delaunay Link Area:  $A_d = h_1 l_1$

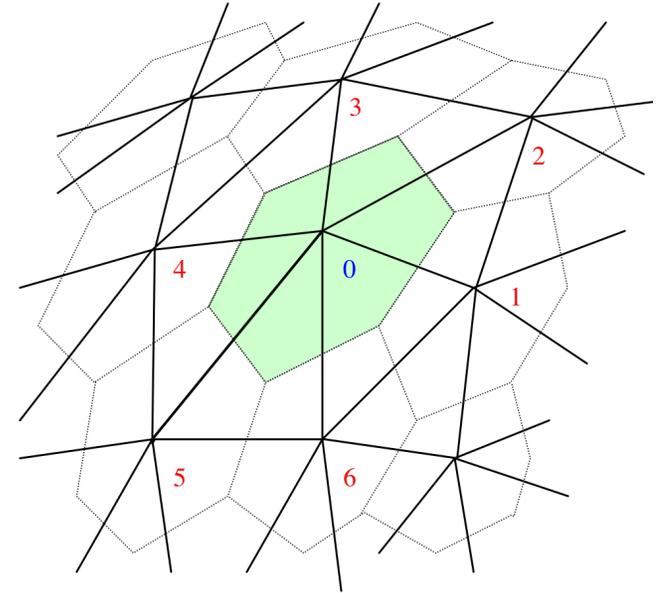
$$\sum_{\Delta_{kij}} \sqrt{g(k)} g^{ij}(k) \frac{(\phi_k - \phi_i)(\phi_k - \phi_j)}{l_{ki} l_{kj}}$$

H. Hamber, S. Liu, *Feynman rules for simplicial gravity*, NP B475 (1996)

# Einstein Regge Curvature

$$\delta_v = 2\pi - \sum_{i \in V} \theta_i$$

$$\sum_v \delta_v = 2\pi\chi = 2\pi(F - E + V)$$



Could Optimize adaptive Delaney triangles on unit sphere

$$\int d^2x \sqrt{g} [\lambda - kR^2 + aR^2] \implies \sum_v A_v [\lambda - 2kR_v + aR_v^2]$$

$$R_v = 2\delta_v / A_v$$

flat triangles:  $\delta_v = 4\pi / A_v$

# FEM have “spectral fidelity”

- Taylor expansion on hypercubic lattice:

$$a^{-1} \sum_{\pm\mu} (\phi(x) - \phi(x + a\mu))^2 \simeq (\nabla\phi)^2 + O(a^2)$$

- Taylor series for FEM does not work!

$$a^2 \sum_y K(x, y) (\phi(x) - \phi(y))^2 \simeq c_{\mu\nu} \partial_\mu \phi \partial_\nu \phi + O(a^2)$$

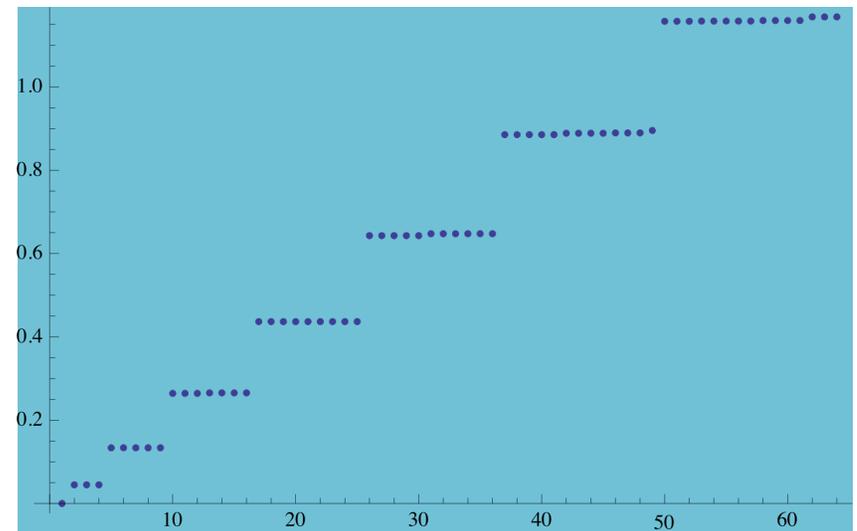
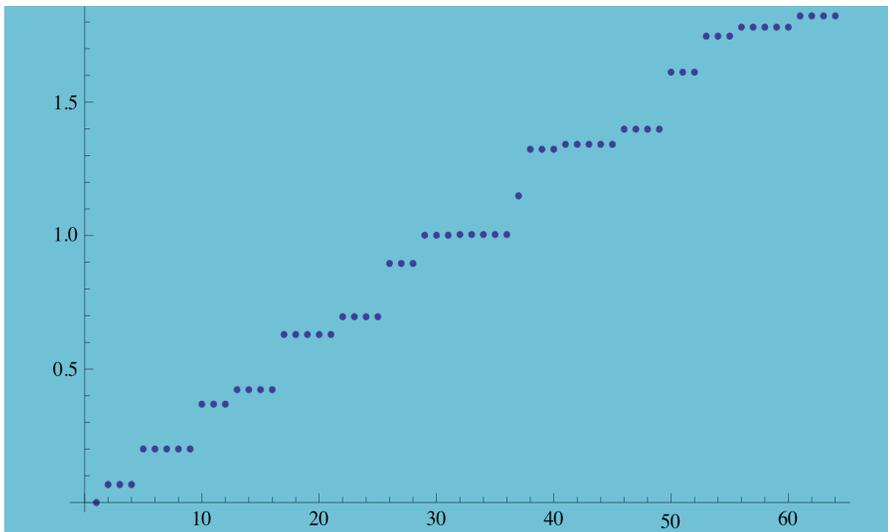
- FEM “meta-theorem”: spectrum < cut-off to  $O(a^2)$  if “triangles are regular enough”

# *FEM fixes the huge Spectral defects*

For  $s = 8$  first  $(l+1)*(l+1) = 64$  ev

BEFORE (K = 1)

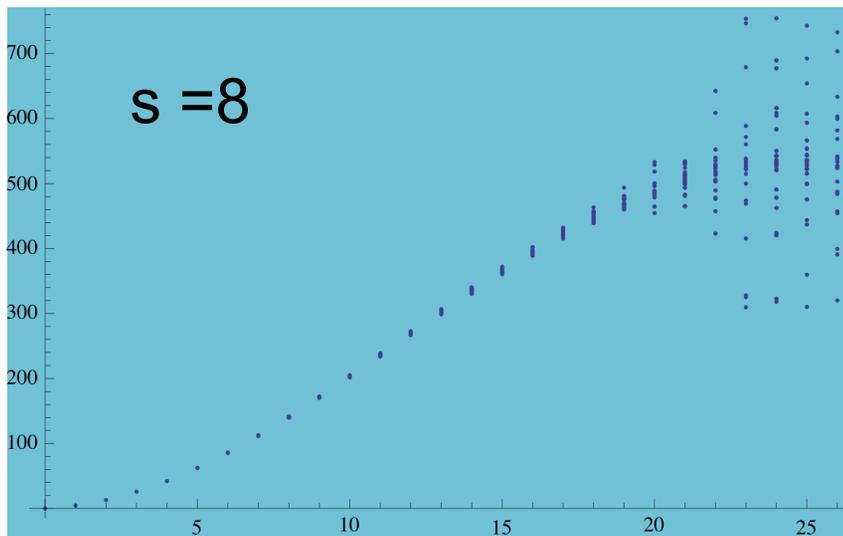
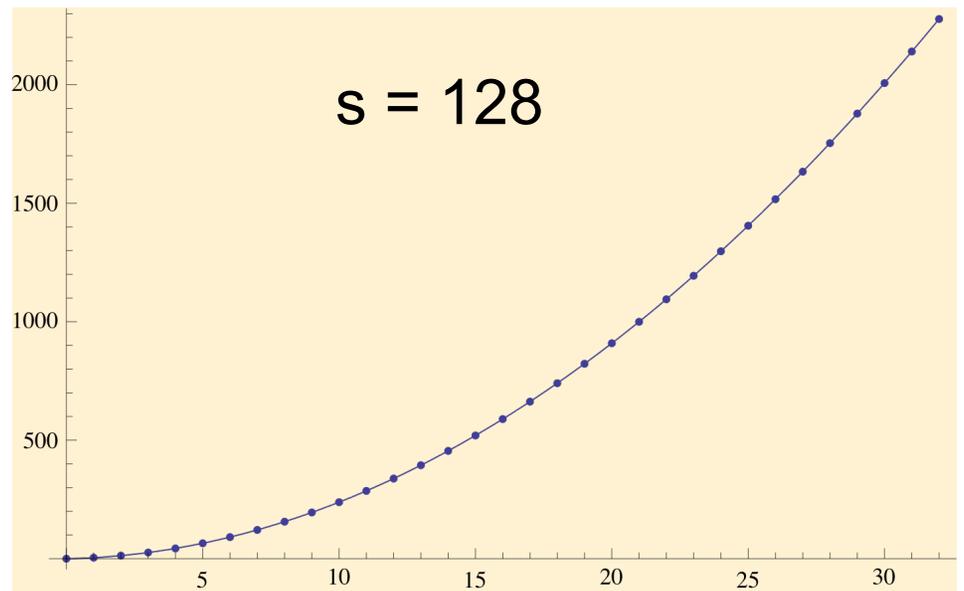
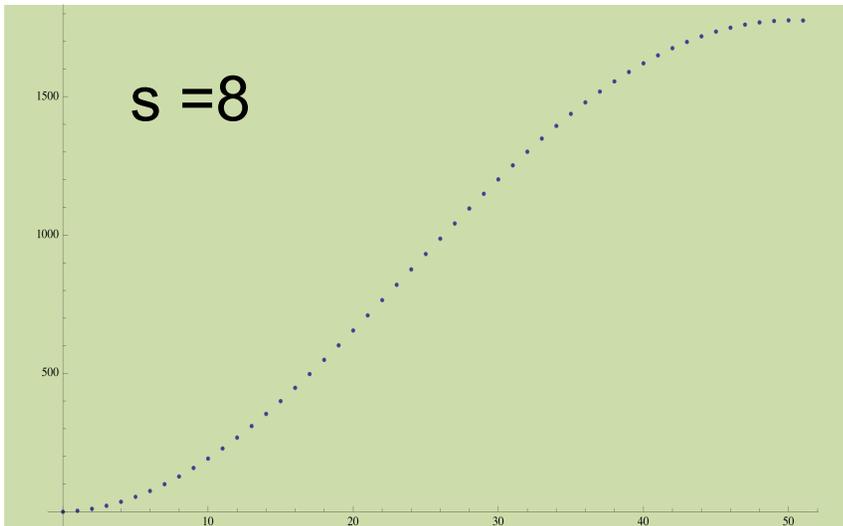
AFTER (FEM K's)



$l, m$

$l, m$

# Spectrum of FE Laplacian on a sphere



Fit

$$l + 1.00012 l^2$$

$$- 13.428110^{-6} l^3 - 5.5724410^{-6} l^4$$



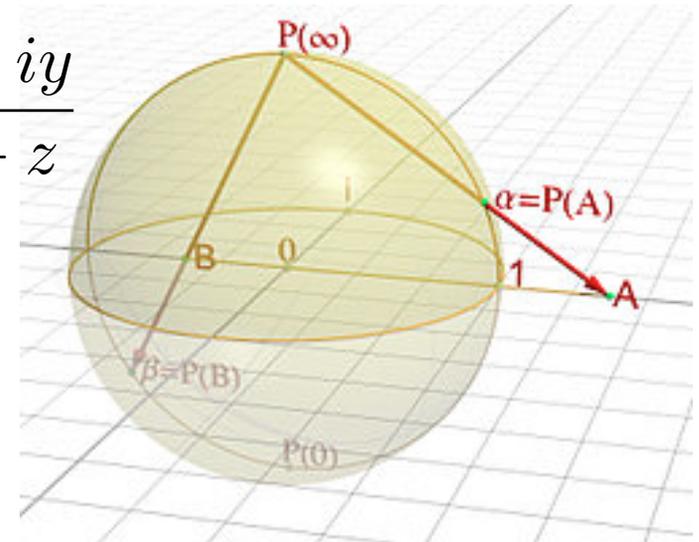
## 2D test: Conformal Projection to Riemann Sphere

projection  $\xi = \tan(\theta/2)e^{-i\phi} = \frac{x + iy}{1 + z}$

### Exact Two point function

$$\langle \phi(x_1)\phi(x_2) \rangle = \frac{1}{|x_1 - x_2|^{2\Delta}} \rightarrow \frac{1}{|1 - \cos\theta_{12}|^\Delta}$$

$$\Delta = \eta/2 = 1/8$$



$$x^2 + y^2 + z^2 = 1$$

### 4 pt function $(x_1, x_2, x_3, x_4) = (0, \xi, 1, \infty)$

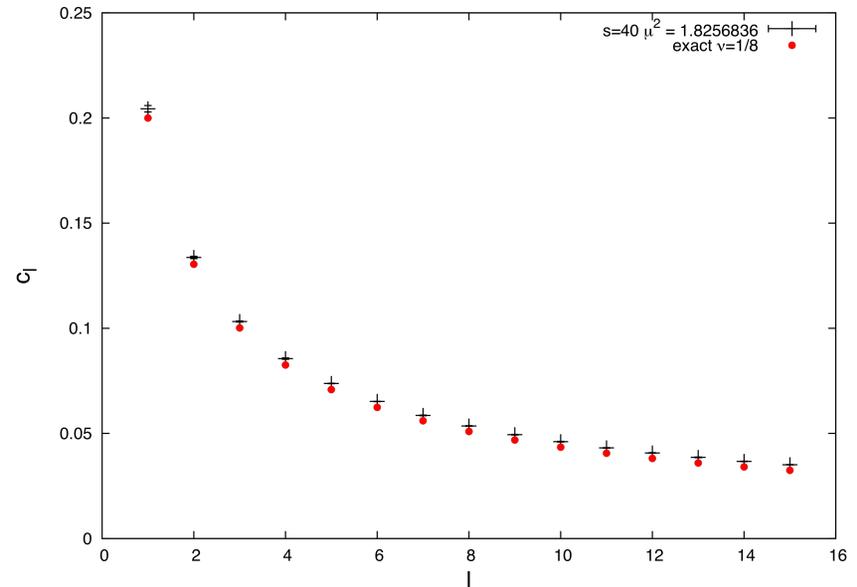
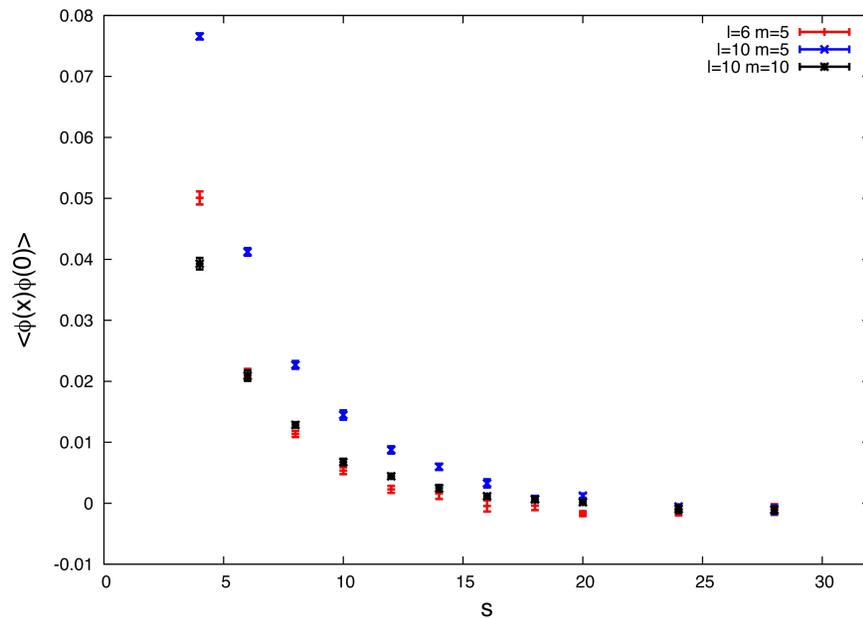
$$g(0, \xi, 1, \infty) = \frac{1}{2|\xi|^{1/4}|1 - \xi|^{1/4}} [1 + \sqrt{1 - \xi} + |1 - \sqrt{1 - \xi}|]$$

### Critical Binder Commulant

$$U_B^* = 1 - \frac{\langle M^4 \rangle}{3\langle M^2 \rangle \langle M^2 \rangle} = 0.567336$$

# Test of rotational symmetry?

Ylm projection.



$$\Delta_\sigma = \eta/2 = 1/8 \simeq 0.128$$

# *The simulation program is running*

(1) Monte Carlo is a “standard” mixture of metropolis, over relax and Wolff methods from:

Ulli Wolff, “Collective Monte Carlo Updating for Spin Systems PRL 62: 361 (1989)

R.C.B. and P. Tamayo, “Embedded Dynamics for  $\phi^4$  Theory”, PRL 62:1087(1989)

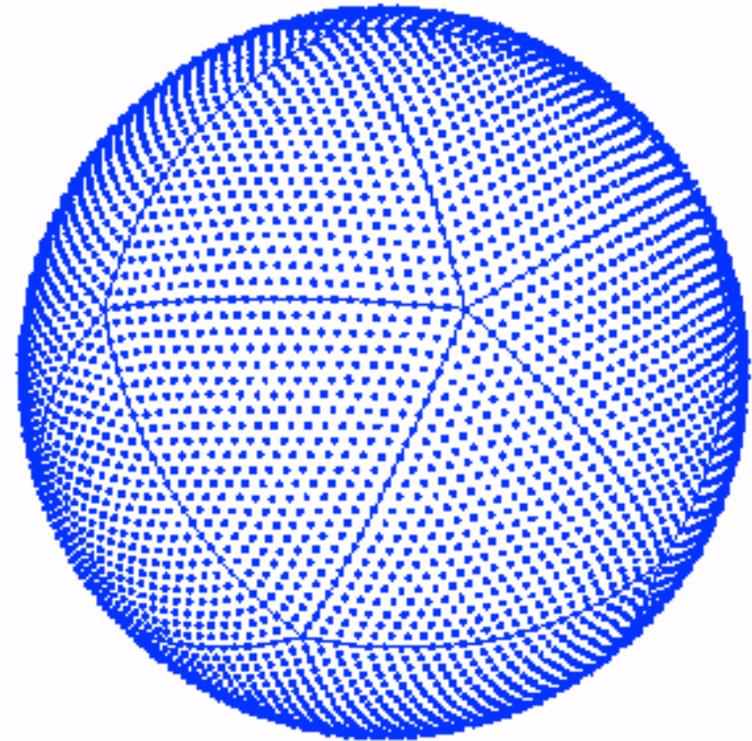
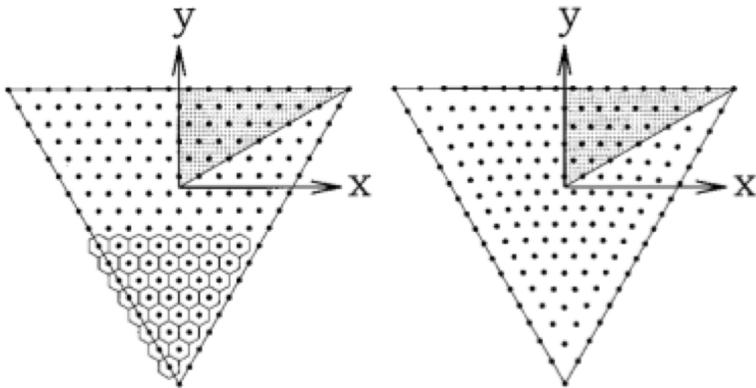
(2) Will compute higher primaries, even Z2 sector, Energy momentum tensor, Conformal Blocks partial waves

(3) The code can run any graph, so we will replace sphere by torus to reproduce  $\phi^4$  numbers from Hasenbusch,...

Hasenbusch, “A Monte Carlo study of leading order scaling corrections of  $\phi^4$  theory on a three-dimensional lattice” J.Phys. A 32 (1999) 4851 \*

# Accurate Results require Better Triangulation, Counter Terms etc

Equal AREA simplex  
for WMAP Ski!



*An icosahedron-based method for pixelizing the celestial sphere*  
MAX TEGMARK THE ASTROPHYSICAL JOURNAL, 470:L81–L84, 1996 October 20

# *Future Challenges & Directions*

- Many extensions are interesting

- Easier problems:

- Prove radial quantization for  $O(N)$  at large  $N$
    - Strengthen bootstrap inequalities for spin systems?

- Harder Problems:

- Gauge fields (with discrete Christoffel connection)?
    - Fermions (with discrete spin connection) ?
    - Flow from UV to conformal IR fixed points for BSM?  
(Cross over from UV, to Hamiltonian to Dilation spectrum)



# EXTRA SLIDES

# Fancy stuff: “FEM” or discrete Exterior Calculus

Need a local reference tangent plane  $\xi^a(x)$  at  $x$ .

Introduce an ortho normal basis in the tangent space:  $\vec{e}_a(x)$

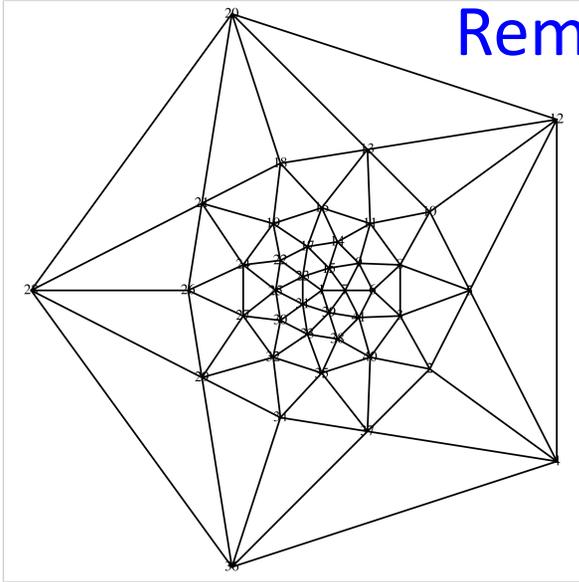
$$g_{\mu\nu}(x) = e_{\mu}^a(x)e_{\nu}^a(x) \quad \bar{\psi}e_a^{\mu}\gamma^a D_{\mu}\psi$$

**Lattice Fermions** are on simplicial complex lattice manifolds with great care! Spin connection has be done carefully.

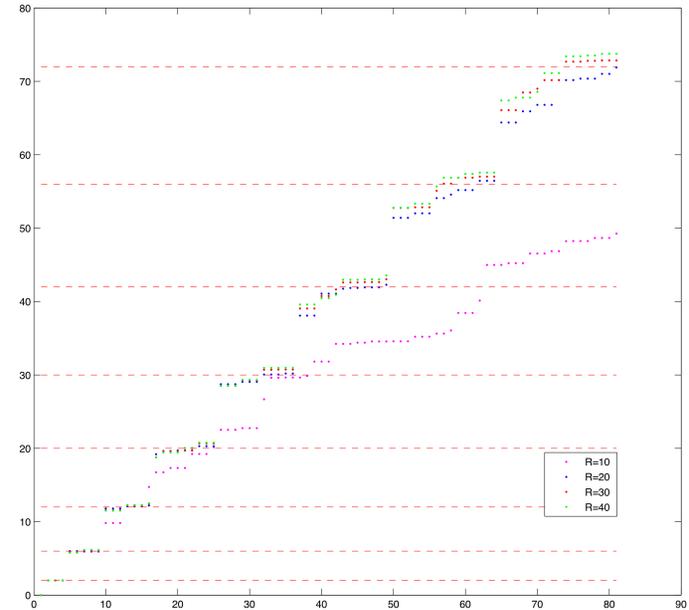
**Compact gauge links** can be represented also a la Christ, Friedberg, Lee! In weak field limit maybe equivalent to using **Nedelic/Whitney “edge” elements** etc.

*\*Simplicial differential form, deRham complex a la Regge Calculus!*

## Remarks on Riemann Sphere



Delaunay triangulation  
commutes with projection!



Spectrum on Cubic Sphere!

- Testing
  - Exact 2D Ising correlators, exponents, central charge, etc.
  - Improved spherical and/, 2<sup>nd</sup> order elements, etc
  - Dynamical  $R^2$  Curvature Constrained Triangulation.